

MESHLESS PETROV-GALERKIN METHOD APPLIED TO AXISYMMETRIC PROBLEMS

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Abstract

An axisymmetric Meshless Local Petrov-Galerkin (MLPG) algorithm is presented for the potential and elasticity problems. In this algorithm the trial and test functions are chosen from different spaces. By a judicious choice of these functions, the integrals involved in the weak form can be restricted to a local neighborhood. This makes the method truly meshless. The MLPG algorithm is used to study various potential and elasticity problems for which exact solutions are available. The sensitivity and effectiveness of the MLPG algorithm to various parameters such as the weight functions, basis functions and support domain radius, etc. was studied. The MLPG algorithm yielded accurate solutions for all weight functions, basis functions and support domain radii considered for all of the problems studied.

Introduction

In the past decade, several meshless methods for the solution of partial differential equations in science and engineering appeared in the literature [1-4]. Meshless methods retain all the advantages of the finite element method without most of the disadvantages such as, element locking and discontinuous derivatives of the secondary variables across the element boundaries. Atluri and Zhu [4] presented a new and innovative meshless approach that uses Petrov-Galerkin weight functions instead of the traditional Galerkin weighted residual method. This meshless method is 'truly' meshless. This method, referred to as the Meshless Local Petrov-Galerkin (MLPG) method, does not need

"mesh" or "cells" either to interpolate the solution variables or to evaluate the integrals that appear in the weak form, but rather needs only nodes that are randomly distributed in the domain.

In the present work, MLPG algorithm is formulated and implemented for axisymmetric potential and elasticity problems. Various choices of the basis functions, forms of weight functions, and sizes of support domain are considered. The approximate solutions obtained by the MLPG algorithm are compared with the exact solutions for evaluating their accuracies.

Axisymmetric Potential Problems

Consider a Poisson's equation for an axisymmetric problem bounded by a toroidal domain with its cross section defined by Ω as shown in Figure 1,

$$\frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{\partial^2 u}{\partial z^2} = g(r, z) \quad \text{in } \Omega \quad (1)$$

with boundary conditions

$$u = \bar{u} \quad \text{on } \Gamma_u \quad \text{and} \quad q = \bar{q} \quad \text{on } \Gamma_q \quad (2)$$

where $\Gamma = \Gamma_u + \Gamma_q$ and $q = du/dn$. In a steady-state heat transfer problem, that is a typical Poisson's problem, u is a temperature function, q is heat flux through the boundary, and g is function of internal heat generation.

The solution for Eq. (1) is sought in a weighted residual manner as

$$2\pi \int_{\Omega} (\nabla^2 u - g) \cdot v \cdot r \cdot dr \cdot dz = 0 \quad (3)$$

where v is a weight function. The factor 2π in Eq. (3) comes from integration over the tangential coordinate θ , because of axisymmetry.

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Axisymmetric Elasticity Problems

The stresses σ_{ij} and body force f_i for an elastic continuum need to satisfy the equations of equilibrium condition

$$\sigma_{ij,j} + f_i = 0 \quad (4)$$

where comma indicates the partial differentiation with respect to the variable that follows the comma.

The boundary conditions on Γ bounding the domain Ω are

$$\mathbf{u} = \bar{\mathbf{u}} \text{ on } \Gamma_u \quad \text{and} \quad \mathbf{t} = \bar{\mathbf{t}} \text{ on } \Gamma_q \quad (5)$$

where \mathbf{u} and \mathbf{t} are displacement and traction vectors, respectively.

The weighted residual statements for Eq. (4) for the axisymmetric case can be written as

$$\begin{aligned} 2\pi \int_{\Omega} \left(\frac{\partial \sigma_r}{\partial r} + \frac{\partial \tau_{rz}}{\partial z} + \frac{1}{r} (\sigma_r - \sigma_{\theta}) + f_r \right) \cdot \\ v_r \cdot r \cdot dr \cdot dz = 0 \\ 2\pi \int_{\Omega} \left(\frac{\partial \tau_{rz}}{\partial r} + \frac{\partial \sigma_z}{\partial z} + \frac{1}{r} \tau_{rz} + f_z \right) \cdot \\ v_z \cdot r \cdot dr \cdot dz = 0 \end{aligned} \quad (6)$$

where σ_r , σ_{θ} , σ_z , and τ_{rz} are stresses in cylindrical coordinates (r, θ, z) and v_r and v_z are weight functions for r and z directions, respectively. In this study, the weight functions v_r and v_z are chosen to be identical.

Since the MLPG algorithm for the potential and the elasticity problems are very similar, the potential problem is utilized to demonstrate the formulation in the following two sections.

Moving Least Square Approximation

In the present meshless method, a trial function, $u(\mathbf{R})$, with moving least squares (MLS) scheme is used to approximate (in cylindrical coordinates (r, z) ; $\mathbf{R} = \hat{\mathbf{i}} \cdot r + \hat{\mathbf{j}} \cdot z$, the radius vector from the origin) the distribution of the potential function u over a number of randomly distributed nodes as shown in Fig. 1. The trial function, $u(\mathbf{R})$, can be expressed in a series form as [1-4]

$$u(\mathbf{R}) = \sum_k \varphi_k(\mathbf{R}) \cdot \hat{u}_k \quad (7)$$

where \hat{u}_k is a fictitious value of u at node k and $\varphi_k(\mathbf{R})$ are the shape functions of node k . (Figure 2 illustrates a least squares approximation to the \hat{u}_k values at each of the randomly distributed nodes for a 1-D case.) The shape function, φ_k , can be written as

$$\varphi_k = \sum_{j=1}^m p_j(\mathbf{R}) \cdot [A(\mathbf{R})^{-1} \mathbf{B}(\mathbf{R})]_{jk} \quad (8)$$

with

$$\mathbf{A}(\mathbf{R}) = \sum_{k=1}^N w_k(\mathbf{R}) \cdot \mathbf{p}^T(\mathbf{R}_k) \cdot \mathbf{p}(\mathbf{R}_k) \quad (9)$$

$$\mathbf{B}(\mathbf{R}) = [w_1(\mathbf{R}) \cdot \mathbf{p}(\mathbf{R}_1), \dots, w_N(\mathbf{R}) \cdot \mathbf{p}(\mathbf{R}_N)] \quad (10)$$

where N is the number of the nodes in the domain of influence (see, for example, Ref. 4), $w_k(\mathbf{R})$ is a weight function defined later, and $\mathbf{p}^T(\mathbf{R})$ are the basis functions and are chosen to be

$$\mathbf{p}^T(\mathbf{R}) = \{1, r, z\} \quad (11a)$$

for linear representation ($m = 3$) and

$$\mathbf{p}^T(\mathbf{R}) = \{1, r, z, r^2, rz, z^2\} \quad (11b)$$

for quadratic representation ($m = 6$).

Note that in Eqs. (9) and (10), the matrices $\mathbf{A}(\mathbf{R})$ and $\mathbf{B}(\mathbf{R})$ are of size $m \times m$ and $m \times N$, respectively, and $w_k(\mathbf{R})$ is the weight function associated with the node k . The weight function $w_k(\mathbf{R})$ is chosen such that it is non-zero over an influence sub-domain surrounding the node k and zero outside of the influence domain [1-4]. In this study, the domain of influence of node k is assumed to be a circle of radius l_k . The radius l_k must be large enough to contain at least m nodes in each direction to prevent the numerical singularity of the \mathbf{A} -matrix in Eq. (9). The weight function is chosen such that it is smooth and equals unity at the center of the sub-domain, and equals zero at the boundary and outside of the sub-domain. In this paper, three spline functions with C^1 , C^2 and C^3 continuity are used as weight functions:

For C^1 :

$$w_k(\mathbf{R}) = \begin{cases} 1 - 3\rho_k^2 + 2\rho_k^3 & 0 \leq \rho_k \leq 1 \\ 0 & \rho_k > 1 \end{cases}, \quad (12a)$$

for C^2

$$w_k(\mathbf{R}) = \begin{cases} 1 - 10\rho_k^3 + 15\rho_k^4 - 6\rho_k^5 & 0 \leq \rho_k \leq 1 \\ 0 & \rho_k > 1 \end{cases}, \quad (12b)$$

for C^3 :

$$w_k(\mathbf{R}) = \begin{cases} 1 - 35\rho_k^4 + 84\rho_k^5 - 70\rho_k^6 + 20\rho_k^7 & 0 \leq \rho_k \leq 1 \\ 0 & \rho_k > 1 \end{cases} \quad (12c)$$

where $\rho_k = d_k / l_k$ is the normalized distance and $d_k = \|\mathbf{R} - \mathbf{R}_k\|$ is the distance between center of the support sub-domain \mathbf{R}_k and the general point \mathbf{R} . The functions are shown in Figure 3a for a 1-D domain. For a 2-D domain, the weight function can be illustrated as a ‘smooth tent’ as shown in Fig. 3b for a C1 continuity spline function.

MLPG Formulation

In the current implementation of MLPG formulation, the weighted residual form of Eq. (3) is used. The essential boundary conditions are included in the weighted residual statement using a penalty method as

$$2\pi \left\{ \int_{\Omega} (\nabla^2 u - g) \cdot v \cdot r \cdot dr \cdot dz - \alpha \cdot \int_{\Gamma_u} (u - \bar{u}) \cdot v \cdot r \cdot d\Gamma \right\} = 0 \quad (13)$$

where α is a penalty parameter which is chosen as a large number (1×10^{16} is used in this study.) Using the divergence theorem, Eq. (13) can be recast (dropping the constant 2π) as

$$\begin{aligned} & - \int_{\Omega} \left[\frac{\partial u}{\partial r} \frac{\partial v}{\partial r} + \frac{\partial u}{\partial z} \frac{\partial v}{\partial z} \right] \cdot r \cdot dr \cdot dz - \int_{\Omega} g \cdot v \cdot r \cdot dr \cdot dz \\ & + \int_{\Gamma} \left[\frac{\partial u}{\partial r} n_r + \frac{\partial u}{\partial z} n_z \right] \cdot r \cdot v \cdot d\Gamma \\ & - \alpha \cdot \int_{\Gamma_u} (u - \bar{u}) \cdot v \cdot r \cdot d\Gamma = 0 \end{aligned} \quad (14)$$

where n_r and n_z are the direction cosines of the normal of the domain boundary along the r - and z -directions, respectively.

Unlike the Galerkin method where the trial and test functions are chosen from same space, the Petrov-Galerkin method uses trial and test function from different spaces. In the present work, the trial functions are chosen as in Eq. (7), and the test functions are chosen so that the evaluation of integrals involved in Eq. (14) for the k^{th} node is confined to a local neighborhood \mathbf{R}_k . This judicious choice of the test function is what makes this method truly meshless. In the present implementation, the test function, $v(\mathbf{R}, \mathbf{R}_k)$, is defined in a similar manner as the weight functions in Eq. (12) by replacing the l_k with l_o i.e.,

$$v(\mathbf{R}, \mathbf{R}_k) = w_k(\mathbf{R}) = f(\rho_k) \quad \begin{matrix} 0 < \rho_k \leq 1 \\ = 0 & \rho_k > 1 \end{matrix}, \quad (15)$$

where $\rho_k = \|\mathbf{R} - \mathbf{R}_k\| / l_o$ is the normalized distance.

For an internal node, l_o is selected to be equal or less than the shortest distance between the node and the domain boundary (for example, nodes p and q in Figure 4). For a node on the domain boundary, the sub-domain is a segment of a circular region (for example, node s in Figure 4).

Substitution of the chosen trial and test functions in the weak form of Eq. (14) leads to

$$[\mathbf{K}] \cdot \{\hat{\mathbf{u}}\} = \{\mathbf{f}\} \quad (16)$$

where

$$\begin{aligned} K_{kj} &= \int_{\Omega_s} \left(\frac{\partial \phi_j(\mathbf{R})}{\partial r} \frac{\partial v(\mathbf{R}, \mathbf{R}_k)}{\partial r} + \frac{\partial \phi_j(\mathbf{R})}{\partial z} \frac{\partial v(\mathbf{R}, \mathbf{R}_k)}{\partial z} \right) \cdot r \cdot dr \cdot dz \\ &+ \alpha \cdot \int_{\Gamma_{su}} \phi_j(\mathbf{R}) \cdot v(\mathbf{R}, \mathbf{R}_k) \cdot r \cdot d\Gamma \\ &- \int_{\Gamma_{su}} \left[\frac{\partial \phi_j(\mathbf{R})}{\partial r} n_r + \frac{\partial \phi_j(\mathbf{R})}{\partial z} n_z \right] \cdot v(\mathbf{R}, \mathbf{R}_k) \cdot r \cdot d\Gamma \end{aligned} \quad (17)$$

and

$$\begin{aligned} f_k &= \int_{\Gamma_{sq}} \bar{q}(\mathbf{R}) \cdot v(\mathbf{R}, \mathbf{R}_k) \cdot r \cdot d\Gamma \\ &+ \alpha \cdot \int_{\Gamma_{su}} \bar{u}(\mathbf{R}) \cdot v(\mathbf{R}, \mathbf{R}_k) \cdot r \cdot d\Gamma \\ &- \int_{\Omega_s} g(\mathbf{R}) \cdot v(\mathbf{R}, \mathbf{R}_k) \cdot r \cdot dr \cdot dz \end{aligned} \quad (18)$$

where $\bar{q}(\mathbf{R})$ is the prescribed value of normal flux on Γ_{sq} . In Eqs. (17) and (18), the region Ω_s is the support domain of $v(\mathbf{R}, \mathbf{R}_k)$ and is assumed to be a circle of radius l_o (see Figures 1 and 4). The Γ_{su} and Γ_{sq} are segments of the boundary formed by the intersection of Ω_s and Γ (see Figure 4 nodes s and q). The region Ω_s can be made as small as possible by a judicious choice of the test function, v . Numerical integration is used to evaluate the integrals involved in Eqs. (17) and (18). A 12-point Gaussian quadrature numerical integration is used in the present study.

The axisymmetric elastic problem is formulated on similar lines and hence this formulation is not presented here.

Examples

To evaluate the current MLPG algorithm for potential and elasticity problems, several patch test problems with exact solutions are considered. The sensitivity of the MLPG solution to the basis functions defined in Eq. (11) and the three weight functions $w_k(\mathbf{R})$ selected in Eq. (12) is studied. Since all of the potential and elasticity problems in this study are linear problems, the patch test models are defined with an arbitrary constant a . The various normalized radiuses of domain of influence (l_k/a) and normalized radiuses of support domains (l_o/a) are used to study their influence on the accuracy of the results.

The results of the MLPG method are compared to exact solutions. Two error norms ($\|e_M\|_1$ and $\|e_M\|_2$) are used to evaluate the effectiveness of various parameters. These norms are defined as

$$\|e_M\|_1 = \frac{1}{M} \sqrt{\sum_{j=1}^M (u - u_{exact})^2} \quad (19a)$$

$$\|e_M\|_2 = \frac{1}{M} \sqrt{\sum_{j=1}^M \{(u - u_{exact}) / u_{exact}\}^2} \quad (19b)$$

where M is the total number of randomly distributed internal points in the domain at which the numerical solution is evaluated and compared to the exact solution. Note that these internal points are independent points and are not associated with the nodal points used in the models. A value of $M = 50$ is used in this study.

Potential Problems

Three potential patch tests involving Laplace- ($\nabla^2 u = 0$), and Poisson- ($\nabla^2 u = \text{constant}$) equation problems are considered.

Figure 5 shows two models of the Laplace and Poisson equation problems used for the patch tests. These examples are analyzed using the MLPG algorithm for various prescribed boundary conditions for u and q on the boundaries.

A triangular domain for two mixed boundary value problems for Laplace equation is studied. The exact solution for this patch test problem is $u = r^2 - 2z^2 + 3z$. The radius of the domain of influence l_k is set to equal to $4a$. Two different types of choices for l_o are considered. In the first choice, the distance from each internal node to the closest boundary is computed. These distances are used as the values of l_o for that node. Then the smallest value of all of the l_o values is used for all nodes on the boundaries. With this definition, the support domain radius is different for different nodes. The MLPG algorithm with different values of l_o for each node recovered the exact solution for the Laplace problems. In the second choice, the smallest value of all the l_o values from the first choice is used for all the nodes. Once again, the exact solution was recovered. In these studies, the effects of the radius of the domain of influence, l_k , were also evaluated by varying the radius from $4a$ to $10a$. The accuracy of results are insensitive to the radius of the domain of influence.

The sensitivity of the MLPG solution is studied by further varying the support domain radius l_o . The convergence of the error norm is presented in Figure 6. All of the nodes in the model were set to have the same value of l_o . The results show that, for Laplace equation, all three weight functions gave excellent solutions for both the mixed boundary values problems and for a variety of the normalized support domain radius, l_o/a .

The problem involving a Poisson's equation over an L-shaped domain is considered next. The exact solution for the patch test problem is $u = r^2 + 2z^2 + 3z - 2$. The convergence of the solution for various values of l_o and for different boundary conditions is presented in Figure 7. Once again, the exact solution is recovered by MLPG algorithm for all idealizations demonstrating that the present MLPG algorithm is very efficient and accurate.

Finally, a heat transfer problem involving radial heat flow in a hollow circular cylinder with prescribed

constant temperatures T_1 and T_2 at inner and outer surfaces, respectively, is considered (see Figure 8). The exact solution of the problem is $T = T_1 - (T_1 - T_2) \cdot \frac{\ln(r/r_1)}{\ln(r_2/r_1)}$. The model used in the

MLPG algorithm for the problem with 25 randomly distributed nodes and the boundary conditions used are shown in Figure 8.

The exact solution for the temperature is a natural logarithmic function and since the exact solution cannot be represented by polynomial, a study of use of various basis functions (polynomial function) is undertaken. Three polynomial basis functions, linear, quadratic and cubic are considered. Note that the higher order basis function requires longer computer times. Figure 9 presents the error norm for various basis functions with a C^1 weight function. The support domain radius, l_o , is set equal to $0.35a$ for the nodes on the boundary and equal to the distance to the closest boundary for the internal nodes. Very accurate solutions are obtained for all cases studied. Figure 9 shows that the higher order basis function yielded more accurate results than the lower order basis function.

Elasticity Problems

Several elasticity problems were studied to evaluate the MLPG algorithm. The elastic material properties, Young's modulus $E = 100$ GPa and Poisson's ratio $\nu = 0.3$, were used in these problems. As in the case of the potential problems, the accuracy of the results are insensitive to the radius of the domain of influence, l_k , if the value of l_k greater than $4a$.

First the axisymmetric elastic MLPG method was evaluated on problems involving axial tension and radial compression of a circular cylinder with different boundary conditions. Since exact solutions for these simple patch test problems can be represented by the polynomial functions, the meshless method, as expected, recovered the solutions to machine accuracy for all weight and basis functions all of the models.

Lamé's Cylinder

A hollow cylinder subjected to constant internal (p_1) and external (p_2) pressures (Lamé's Problem) is considered (see Figure 10). For the problem analyzed, internal pressure p_1 , external pressure p_2 , the inner radius $r_1 = a$ and the outer radius $r_2 = 2a$ were used. The models used for the problem are shown in Figure 10. The exact solution of the problem is

$$\sigma_r = C_1 - C_2 / r^2 ; \sigma_\theta = C_1 + C_2 / r^2 ; \sigma_z = 2 \cdot \nu \cdot C_1$$

$$u_r = \frac{1+\nu}{E} \cdot \left[C_1 \cdot (1-2\nu) \cdot r + \frac{C_2}{r} \right] ; u_\theta = u_z = 0 ;$$

where

$$C_1 = \frac{p_1 \cdot r_1^2 - p_2 \cdot r_2^2}{r_2^2 - r_1^2} ; C_2 = \frac{p_1 - p_2}{r_2^2 - r_1^2} \cdot r_1^2 \cdot r_2^2 \quad (20)$$

Figure 11 shows the variation of the error norm for the three different weight functions with the normalized domain size using quadratic basis function. All of the nodes in the model were assumed to have the same support domain radius, l_o . The error norm reduces substantially as the domain size increases for all the three weight functions considered. Figure 11 also shows that the error norm of C^1 -weight functions is lower than those with C^2 - and C^3 - weight functions. One of the reasons why C^1 -weight function produces less error than the C^3 -weight function is possibly due to the fact that C^1 function has a larger "effective radius" than the C^3 function. The effective radius is defined to be a value of ρ_k in Eq. (12) as the value of the corresponding weight function $w_k(\mathbf{R})$ is too small to operative in the integrals in Eq. (17). For an arbitrary small threshold value of the weight function, $w_k(\mathbf{R}) = 0.01$, the effective radius of C^1 function is 0.94 while the C^3 function is about 0.86 (see Figure 3(a)). This suggests that for Lamé's problem the effective radius appears to be influencing the accuracy of the results.

An additional study of the domain size effect is conducted using a model with 25 randomly distributed nodes (Figure 12). Two different types of choices for l_o are considered namely equal and unequal size domain. In the first choice "equal size domain", the l_o is kept the same for all of the nodes and equals to the smallest distance from an internal node to the closest boundary (see Figure 12). In the second choice, "unequal size domain" is used. In this choice, for each internal node the l_o is kept as large as possible and set to be the distance of the node to the closest boundary. For each boundary node, the l_o is kept to the distance from that node to its adjacent boundary nodes (see Figure 12). The results in Figure 11 suggest that a more accurate solution (smaller error norm) is obtained when a larger domain size is used. Thus, the unequal size domain is expected to yield a more accurate solution. This expectation is confirmed in Figure 12 with the randomly distributed node model.

The effects of the basis function are also evaluated, since the exact solution for the displacement in Lamé's problem can not be represented precisely by any of the three – linear, quadratic and cubic - polynomial basis functions used. The error norm obtained with each of the choices of the basis functions is presented in Figure 13. The cubic basis function yielded the most accurate

results while the linear basis function produced a marginally better result than the quadratic function.

Rotating Disk

The next problem studied is that of a disk with uniform thickness rotating with a speed of ω rad/sec. as shown in Figure 14. In this problem, all the boundaries are stress free while the disk is subjected to an inertial force per unit volume that equals to $\rho \cdot \omega^2 \cdot r$, where ρ is the density of the disk material and ω is the angular velocity. The exact solution of the problem when the Poisson's ratio is equal to zero is given by [see pages 335-337, Ref. 5]

$$\begin{aligned}\sigma_r &= \frac{3}{8} \cdot \rho \omega^2 (r_1^2 + r_2^2 - \frac{r_1^2 \cdot r_2^2}{r^2} - r^2) \\ \sigma_\theta &= \frac{3}{8} \cdot \rho \omega^2 (r_1^2 + r_2^2 + \frac{r_1^2 \cdot r_2^2}{r^2} - \frac{1}{3} r^2) \\ \sigma_z &= 0; \quad u_z = 0 \\ u_r &= \frac{3}{8} \cdot \rho \omega^2 \left[r \cdot (r_1^2 + r_2^2) + \frac{r_1^2 \cdot r_2^2}{r} - \frac{1}{3} r^3 \right] \quad (21)\end{aligned}$$

where again E is Young's modulus of the material.

Three models are used to study the convergence of the solution to nodal refinement (Figure 14). In the model-1, seven nodal points are placed in the r direction. In the model-2, the distance between adjacent nodes is half of the distance between adjacent nodes in model-1 ($d=\Delta/2$.) In the model-3, the distance between adjacent nodes is a quarter of the distance in model-1 ($d=\Delta/4$.) The number of nodes in z direction is kept the same for all of three models.

Figure 15 shows the effect of the nodal arrangement with constant domain size (equals to the distance between the nodes of model-3, i.e. $l_o = \Delta/4$). The model-3 has most nodal points and produced, as expected, the most accurate solution. However, the model-3 requires about 4 times the CPU time of model-1. The effect of the basis functions is shown in Figure 16 using model-2 with C^1 weight function. The cubic basis function yielded a more accurate result than both of the linear and the quadratic basis functions. However, the error norm of the linear basis function accuracy is about 0.1%.

Figure 17 shows the effect of the domain size using model-2 and a cubic basis function. The results, once again, show that the large support domain radius yields a small error norm.

Concluding Remarks

A meshless Local Petrov-Galerkin (MLPG) algorithm is presented for analyzing potential and elasticity problems in axisymmetric domains. In this method, the trial and test functions are chosen from different spaces. By a judicious choice of these functions, the integrals involved in the weak form can be restricted to a local neighborhood. This makes the method truly meshless. The MLPG algorithm is implemented and its efficiency is studied with respect to three basis functions, three different weight functions, and different sizes of local support domain.

Potential patch test problems involving Laplace and Poisson equations are used to evaluate the efficiencies of the algorithm. To evaluate the effectiveness of the MLPG algorithm for elasticity problems, several simple patch test problems, a Lamé's cylinder problem and rotating disk problem, are studied.

The C^1 weight function yielded more accurate results than C^2 and C^3 functions for the patch test problems studied. Since the C^1 function requires lesser computing time than C^2 and C^3 functions, it should be used in most problems. The larger value of the support domain radius, l_o , tended to decrease the error norm. The variation of the radius of the domain of influence, l_k , does not show a significant effect on the accuracy of the result. For most problems studied, the cubic basis function yielded most accurate results. Increasing the order of the basis function substantially increases the computing time, but reduces the error norm.

In summary, the current MLPG algorithm yielded accurate solutions for all the potential and elasticity patch tests and other problems studied. A simple C^1 weight function, with a large support domain radius and a simple linear basis function are recommended for accurate results.

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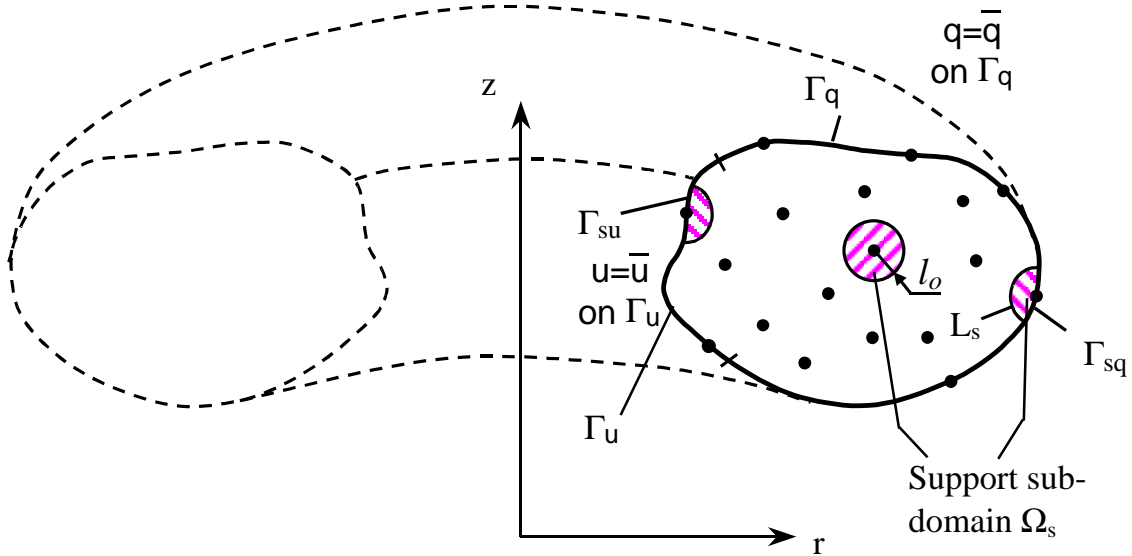


Figure 1 A typical axisymmetric potential problem

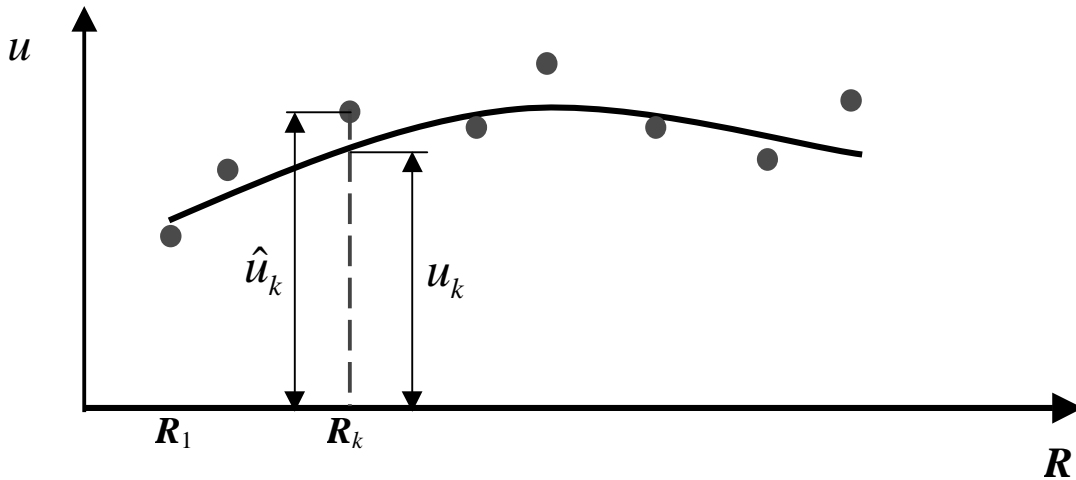
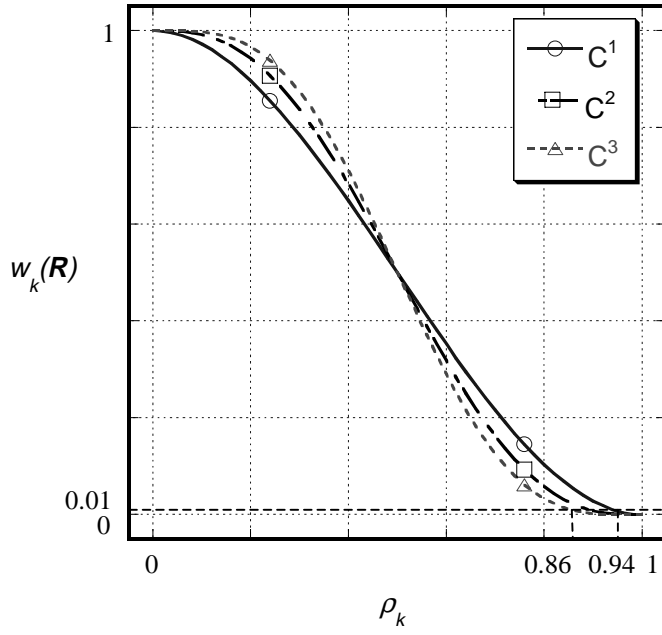
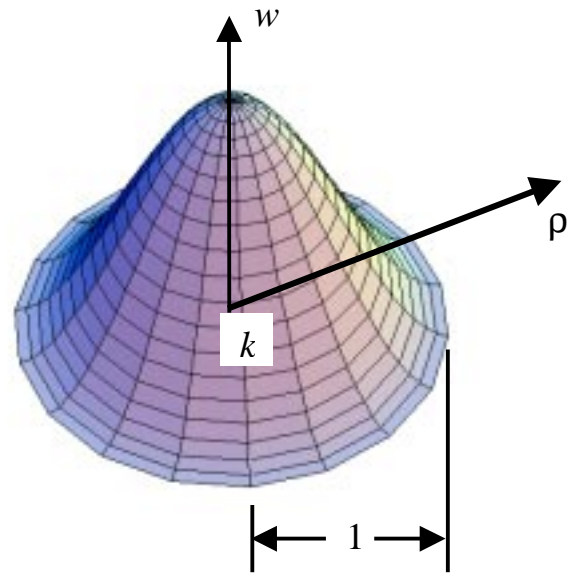


Figure 2 Approximation for u using MLS scheme for a 1-D case



(a) Various spline weight functions used (1-D representation)



(b) C^1 spline weight function (2-D representation)

Figure 3 Spline weight functions used

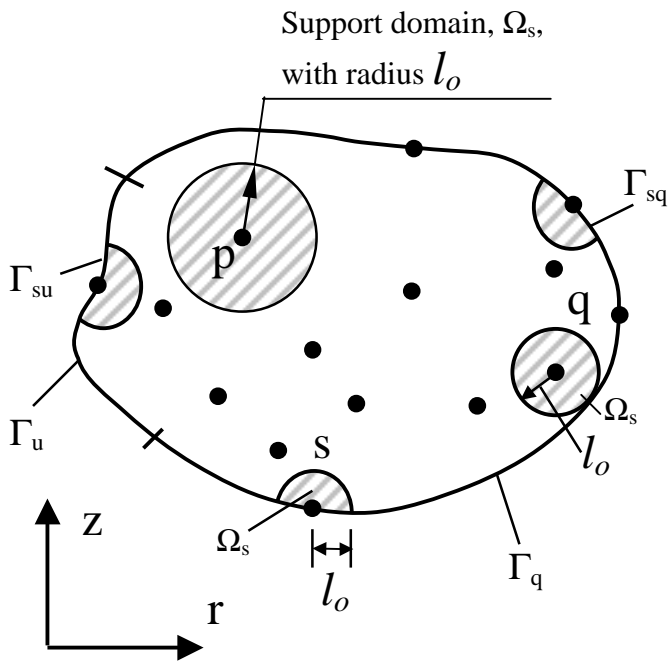
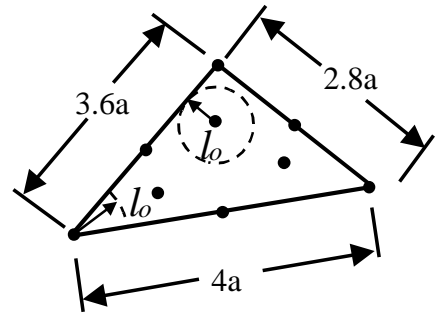
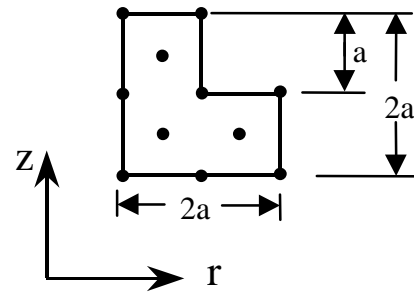


Figure 4 Convenient definitions of l_o in the MLPG method

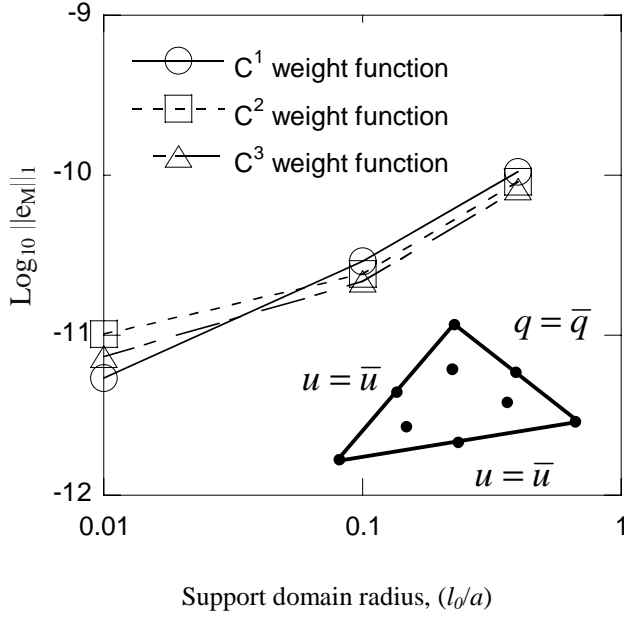


(a) Triangular domain

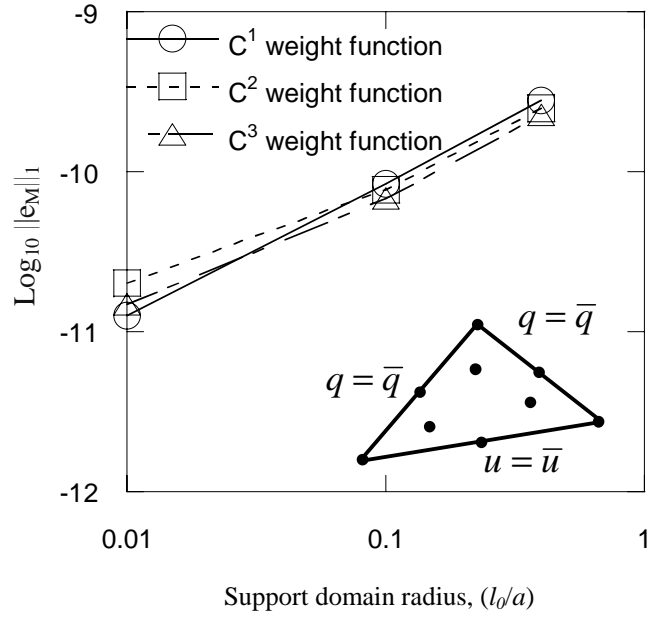


(b) L-shaped domain

Figure 5 Models used in the patch tests of Laplace and Poisson's Equations

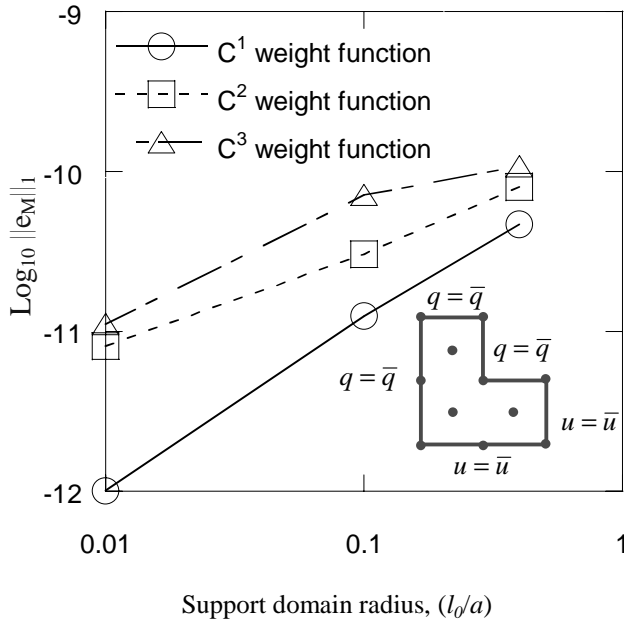


(a) Mixed boundary value problem –1

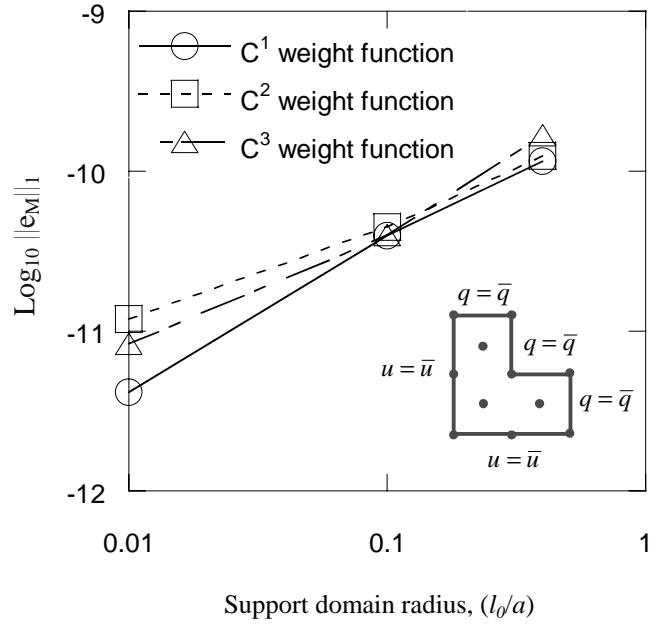


(b) Mixed boundary value problem –2

Figure 6 Laplace problem involving a triangular domain - Convergence of the solution for various values of l_0 . ($\nabla^2 u = 0$, $u_{exact} = r^2 - 2z^2 + 3z$; Quadratic basis; $l_k = 4a$)



(a) Mixed boundary value problem –1



(b) Mixed boundary value problem –2

Figure 7 Poisson's problem involving an L-shaped domain - Convergence of the solution for various values of l_0 . ($\nabla^2 u = 8$, $u_{exact} = r^2 + 2z^2 + 3z - 2$; Quadratic basis; $l_k = 4a$).

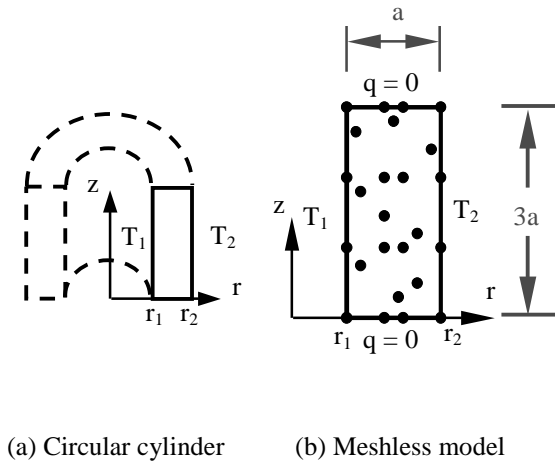


Figure 8 Heat transfer problem in a hollow circular cylinder

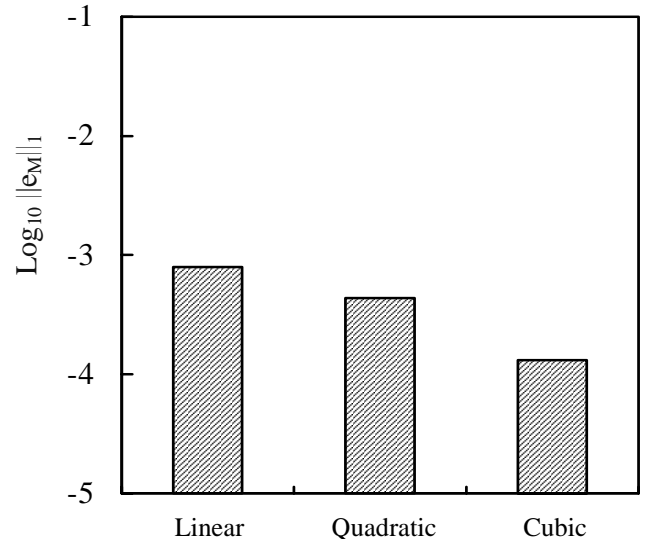


Figure 9 Heat transfer problem involving a hollow circular cylinder – Effect of basis functions.

$$(T_{\text{exact}} = T_1 - (T_1 - T_2) \cdot \frac{\ln(r/r_1)}{\ln(r_2/r_1)}, C^1 \text{ weight function})$$

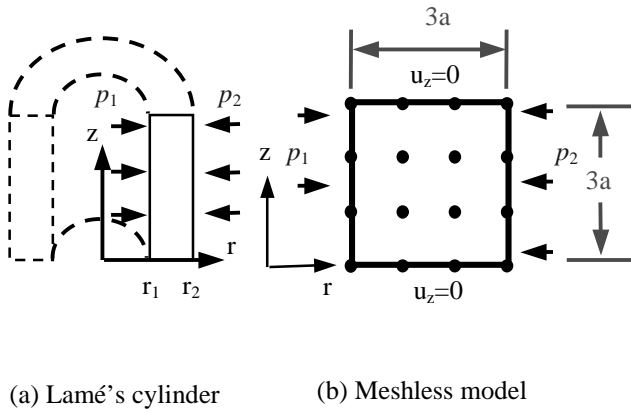


Figure 10 Hollow cylinder subjected to internal and external pressures (Lamé's Problem)

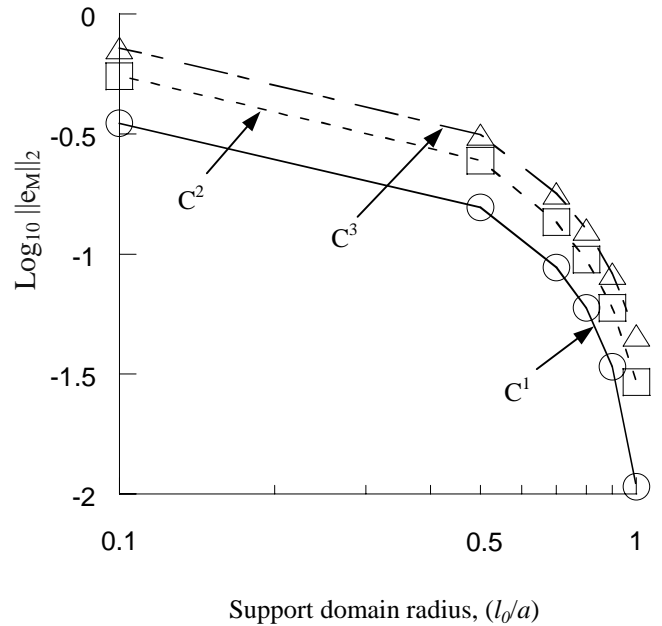


Figure 11 Lamé's cylinder - Effects of domain radius and the weight functions

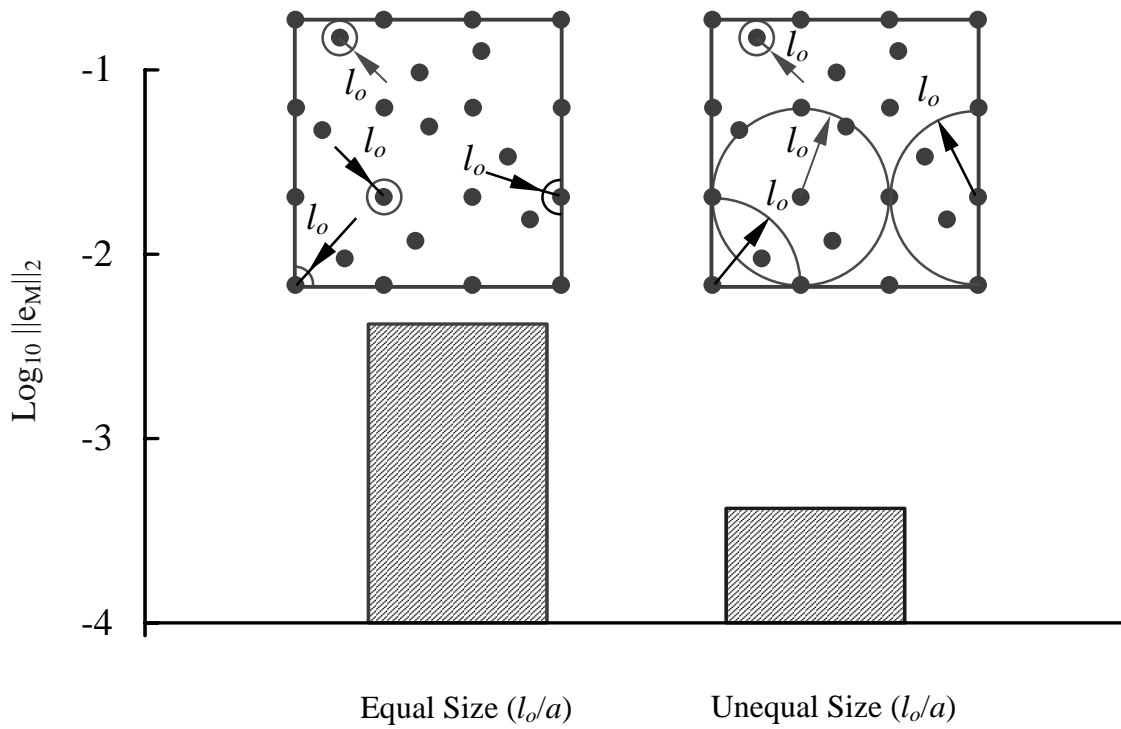


Figure 12 Lamé's cylinder - Effects of support domain radius (l_0/a) - (Cubic basis function, C^1 weight function)

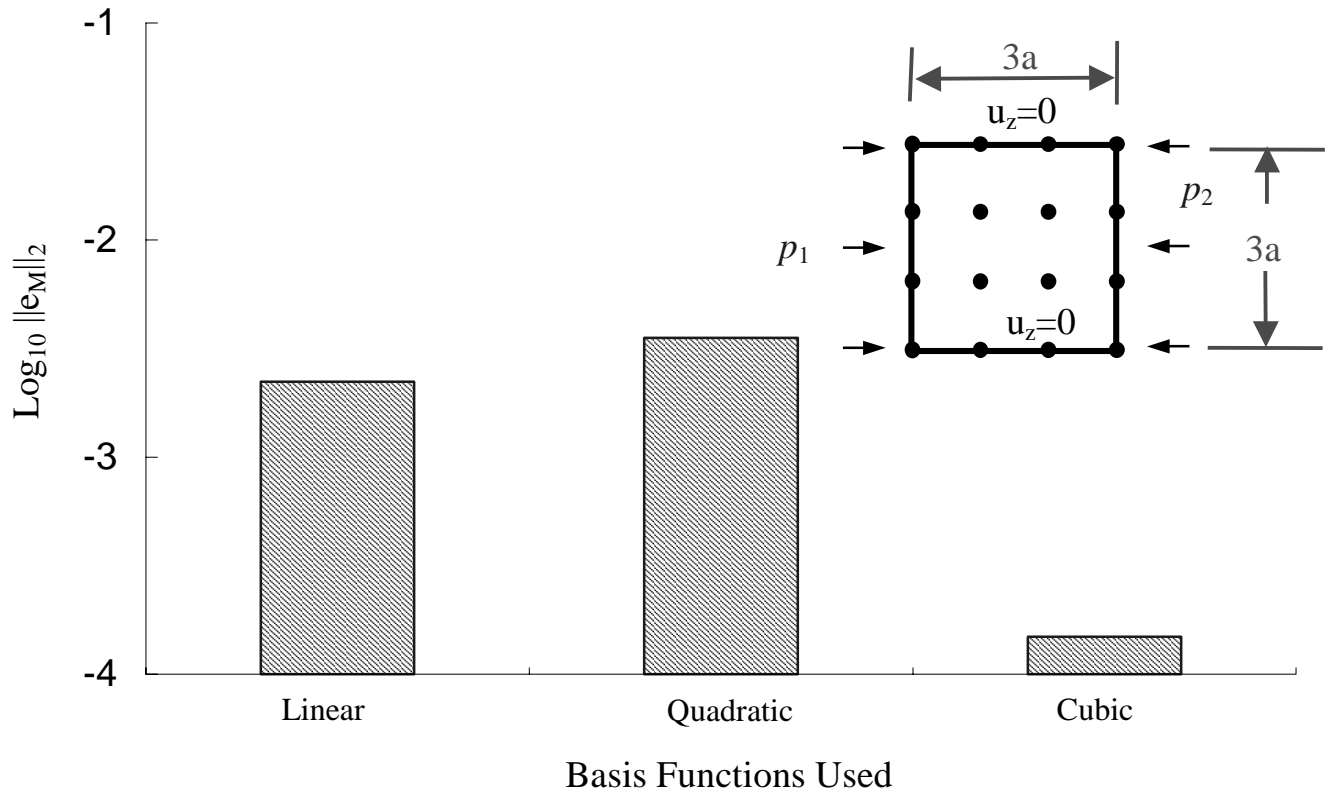


Figure 13 Lamé's cylinder - Effects of basis functions (C^1 weight function; $l_0/a = 0.75$)

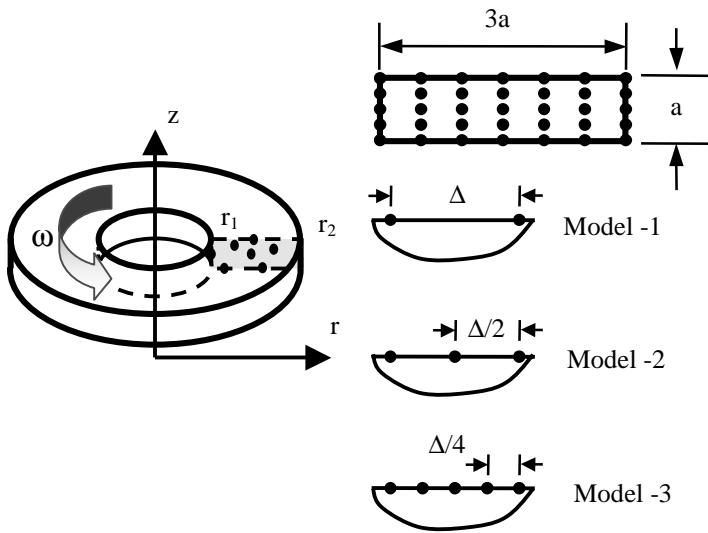


Figure 14 Rotating disk problem and the three models used

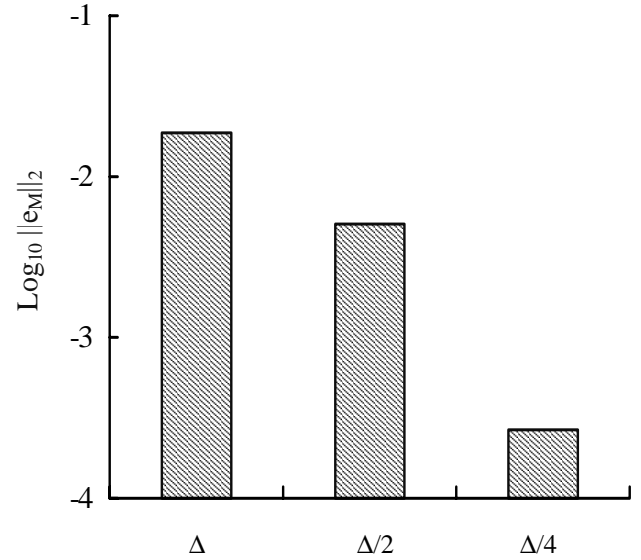


Figure 15 Rotating disk - Effect of nodal refinement - (C^1 weight function; Cubic basis function; $l_o/a = 0.125$)

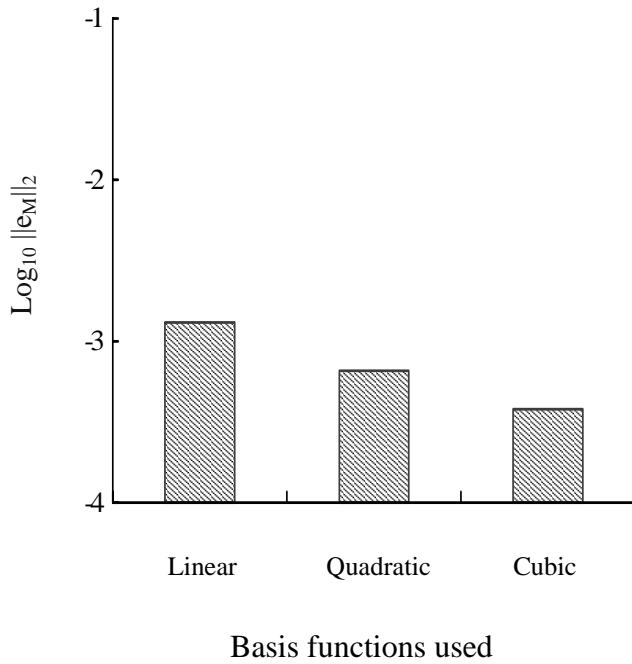


Figure 16 Rotating disk - Effect of basis functions - (C^1 weight function; $l_o/a = 0.25$; model-2)

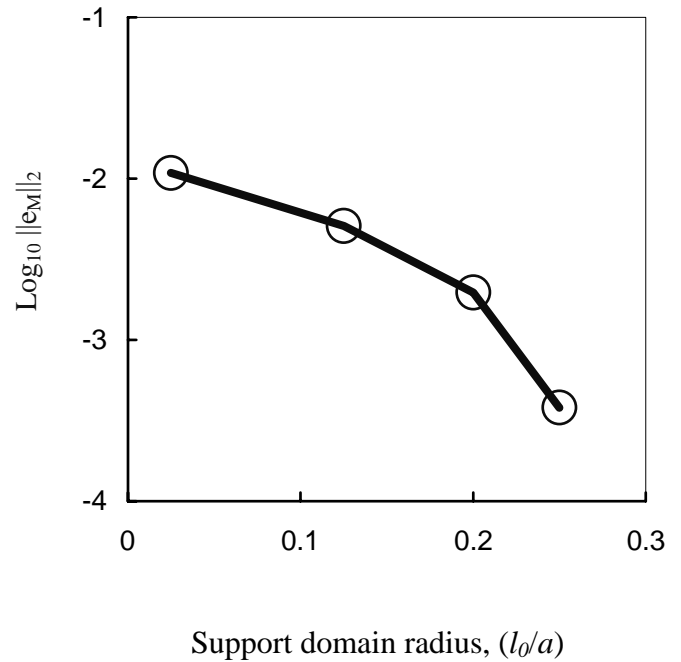


Figure 17 Rotating disk - Effect of support domain radius - (C^1 weight function; Cubic basis function, model-2)